# OPTICS PHYS 311

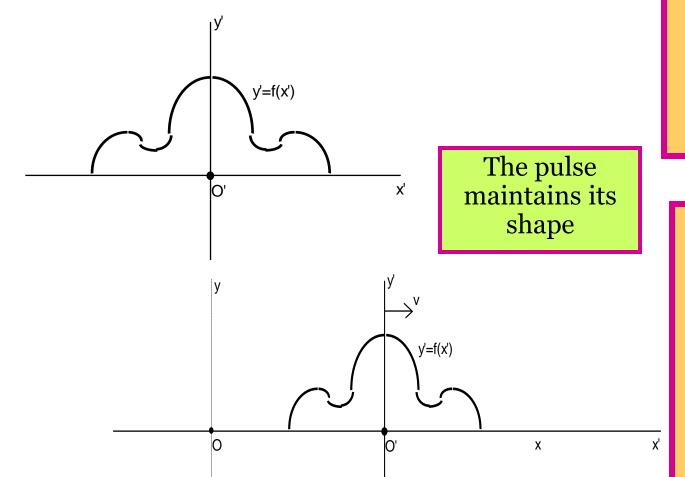
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# WAVE EQUATION

# Introduction

Before we can understand how light moves from one medium to another and how it interacts with lenses and mirrors, we must be able to describe its motion mathematically

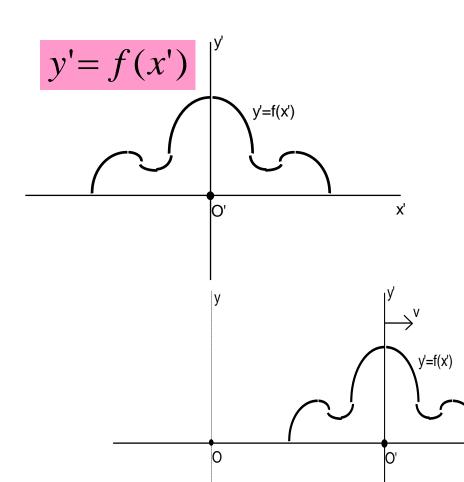
- 1. The most general form of a traveling wave.
- 2. The differential equation it satisfies.



One dimensional wave pulse of arbitrary shape fixed to a coordinate system O'(x',y')

The O' system &
the pulse
move to the
right along
the x-axis at a
uniform speed
v relative to a
fixed
coordinate
system O(x,y)

Optics 311 - Wave Motion



Any point P on the pulse can be described by either of the two coordinates x or x'

$$x' = x - \upsilon t$$

$$y = y' = f(x') = f(x - \upsilon t)$$

Left movement → +v

$$y = f(x \pm \upsilon t)$$



$$\psi = \psi(x, y, z, t) = f(\vec{r} \pm \vec{\upsilon}t)$$

The general form of a traveling wave

$$|\psi(\vec{r},t)|_{t=0} = \psi(\vec{r})$$

Shape or Profile

# **One Dimensional Waves**

$$\psi(\vec{r},t) \rightarrow \psi(x,t) = \psi(x')$$

The variation of ψ with respect to position is given by:

$$\frac{\partial \psi(x,t)}{\partial x} = \frac{\partial x'}{\partial x} \frac{d\psi(x')}{dx'}$$
$$= \frac{d\psi(x')}{dx'}$$

The variation of ψ with respect to time is given by:

$$\frac{\partial \psi(x,t)}{\partial t} = \frac{\partial x'}{\partial t} \frac{d\psi(x')}{dx'}$$
$$= \pm \upsilon \frac{d\psi(x')}{dx'}$$

Taking the second derivative of these yields:

# **One Dimensional Waves**

#### Taking the second derivative of these yields:

$$\frac{\partial^2 \psi(x,t)}{\partial x^2} = \frac{\partial}{\partial x} \left( \frac{\partial \psi(x,t)}{\partial x} \right)$$

$$= \frac{\partial}{\partial x} \left( \frac{\partial \psi(x')}{\partial x'} \right)$$

$$= \left( \frac{\partial x'}{\partial x} \frac{\partial}{\partial x'} \right) \frac{d\psi(x')}{dx'}$$

$$= \frac{d^2 \psi(x')}{dx'^2}$$

$$\frac{\partial^{2} \psi(x,t)}{\partial t^{2}} = \frac{\partial}{\partial t} \left( \frac{\partial \psi(x,t)}{\partial t} \right)$$

$$= \frac{\partial}{\partial t} \left( \pm \upsilon \frac{d \psi(x')}{d x'} \right)$$

$$= \left( \pm \upsilon \frac{\partial}{\partial x'} \right) \left( \pm \upsilon \frac{d \psi(x')}{d x'} \right)$$

$$= \upsilon^{2} \frac{d^{2} \psi(x')}{d x'^{2}}$$

$$\frac{\partial^2 \psi(x,t)}{\partial x^2} = \frac{1}{v^2} \frac{\partial^2 \psi(x,t)}{\partial t^2}$$
THE WAVE EQUATION

#### **Three Dimensional Waves**

#### **Extending to three dimensions**

$$\vec{\nabla} = \frac{\partial}{\partial x}\hat{i} + \frac{\partial}{\partial y}\hat{j} + \frac{\partial}{\partial z}\hat{k}$$

Directional derivative ≡ gradient

#### THE WAVE EQUATION

$$\nabla^2 \psi(\vec{r}, t) = \frac{1}{\upsilon^2} \frac{\partial^2 \psi(\vec{r}, t)}{\partial t^2}$$

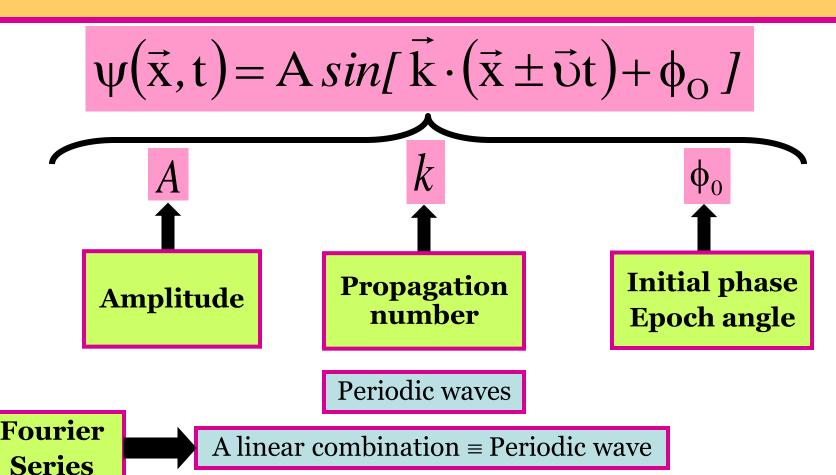
∇² is called the Laplacian operator

$$\nabla^2 = \vec{\nabla} \cdot \vec{\nabla} = \left(\frac{\partial}{\partial x}\hat{i} + \frac{\partial}{\partial y}\hat{j} + \frac{\partial}{\partial z}\hat{k}\right) \cdot \left(\frac{\partial}{\partial x}\hat{i} + \frac{\partial}{\partial y}\hat{j} + \frac{\partial}{\partial z}\hat{k}\right) = \frac{\partial^2}{\partial x^2} + \frac{\partial^2}{\partial y^2} + \frac{\partial^2}{\partial z^2}$$

#### CARTESIAN COORDINATES

# **Harmonic Waves**

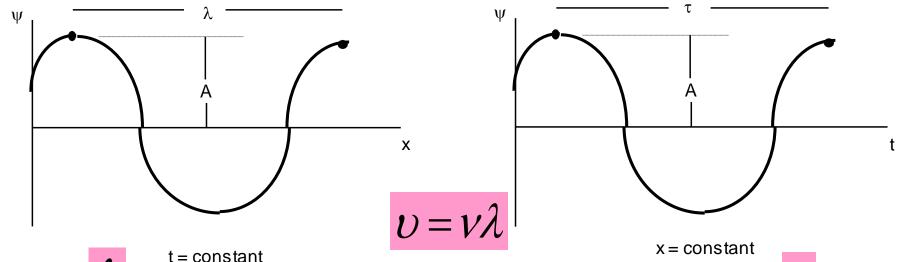
Of special importance are simple harmonic waves that involve the sine or cosine functions



# **Harmonic Waves**

#### What is the physical interpretation of HW equation?

$$\psi(\vec{x},t) = A \sin[\vec{k} \cdot (\vec{x} \pm \vec{\upsilon}t) + \varepsilon]$$



λ

wavelength

$$k = 2\pi/\lambda$$

$$\kappa = \frac{1}{\lambda}$$

Wave number

$$\omega = 2\pi v$$

**Angular frequency** 

$$\tau = \frac{1}{V}$$

# **Harmonic Waves**

$$\psi(\vec{x},t) = A \sin[\vec{k} \cdot (\vec{x} \pm \vec{v}t) + \phi_0]$$

$$\varphi = \vec{k} \cdot (\vec{x} \pm \vec{v}t) + \phi_0$$

**Phase** 

When x and t change together in such a way that  $\varphi$  is constant, the displacement  $\psi = A \sin \varphi$  is also a constant.



Describes the motion of a fixed point on the wave form

$$d\varphi = 0 = \vec{k} \cdot (d\vec{x} \pm \vec{\upsilon}dt) \longrightarrow \frac{d\vec{x}}{dt} = \pm \vec{\upsilon} \longleftarrow \frac{\mathbf{Wave velocity}}{\mathbf{Phase velocity}}$$

 $\phi_0$ 

Initial phase

when 
$$x = 0 \& t = 0$$

$$\psi_{o} = A \sin \phi_{0}$$

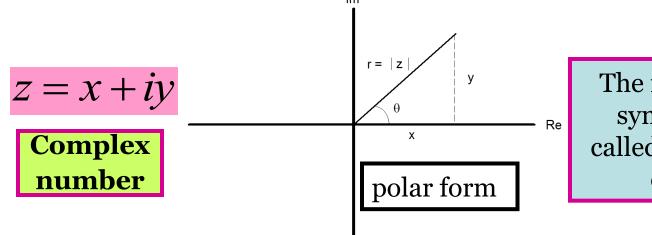
# **Complex Representation of Waves**

# It is convenient to represent wavefunctions as complex functions

$$\psi(\vec{x},t) = A \sin(\vec{k} \cdot (\vec{x} \pm \vec{v}t) + \phi_0)$$

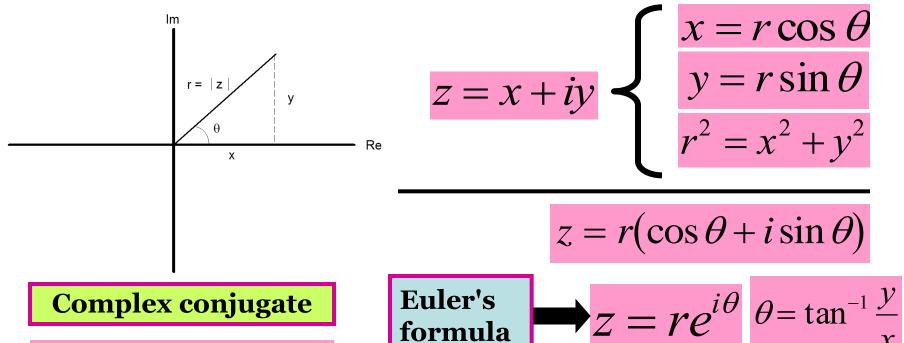
$$\psi(x,t) = Ae^{i(\vec{k}\cdot\vec{x} - \omega t + \phi_0)}$$

mathematically simpler to use



The **magnitude** of *z*, symbolized by *r*, is called **absolute value** or **modulus** 

# **Complex Representation of Waves**



Complex conjugate

$$z^* = x - iy$$
 or  $z^* = re^{-i\theta}$ 

$$\Rightarrow z = re^{i\theta}$$

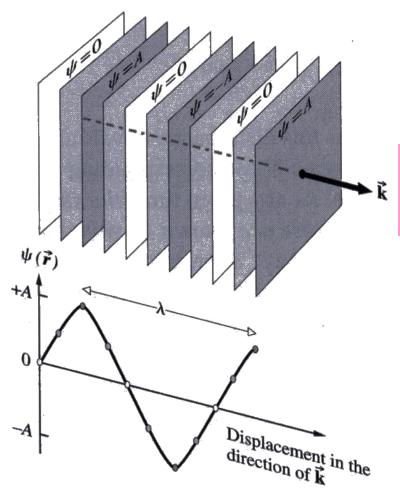
$$\theta = \tan^{-1} \frac{y}{x}$$

$$\psi(x,t) = Ae^{i(\vec{k}\cdot\vec{x}-\omega t+\varepsilon)}$$

$$\operatorname{Re}(\psi) = A\cos(\vec{k}\cdot\vec{x} - \omega t + \varepsilon)$$

$$Re(\psi) = A\cos(\vec{k} \cdot \vec{x} - \omega t + \varepsilon)$$
$$Im(\psi) = A\sin(\vec{k} \cdot \vec{x} - \omega t + \varepsilon)$$

# **Plane Waves**



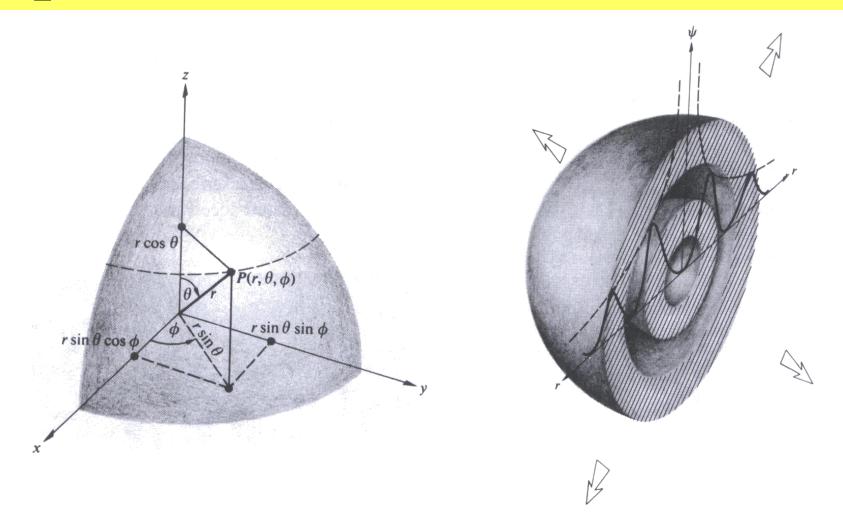
These planes are called the **wavefronts** of the disturbance

$$\nabla^2 \psi = \frac{\partial^2 \psi}{\partial x^2} + \frac{\partial^2 \psi}{\partial y^2} + \frac{\partial^2 \psi}{\partial z^2} = \frac{1}{v^2} \frac{\partial^2 \psi}{\partial z^2}$$

A wave which satisfies the above wave equation is called a **PLANE WAVE** 

$$\psi = Ae^{i\left(\vec{k}\cdot\vec{r}-\omega t\right)}$$

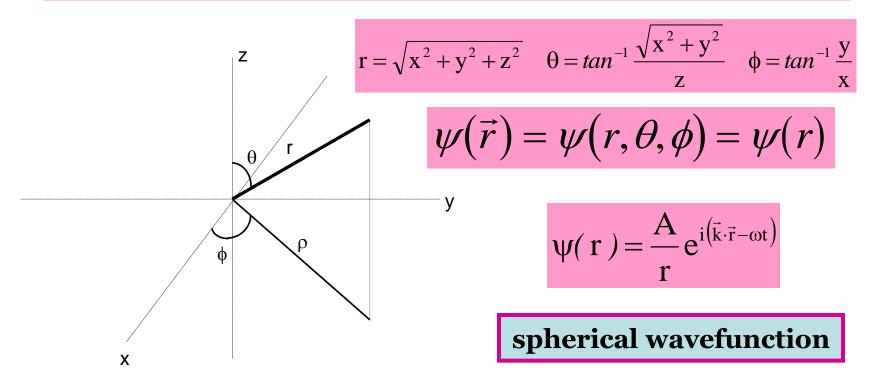
# **Spherical Waves**



# **Spherical Waves**

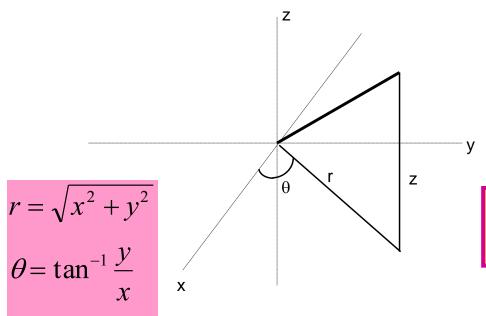
In this coordinate system, the Laplacian operator becomes

$$\nabla^{2} = \frac{1}{r^{2}} \frac{\partial}{\partial r} \left( r^{2} \frac{\partial}{\partial r} \right) + \frac{1}{r^{2} \sin \theta} \frac{\partial}{\partial \theta} \left( \sin \theta \frac{\partial}{\partial \theta} \right) + \frac{1}{r^{2} \sin^{2} \theta} \frac{\partial^{2}}{\partial \phi^{2}}$$



the amplitude of a spherical wave decreases as it moves away from its source

# Cylindrical Waves



In this coordinate system, the Laplacian operator becomes

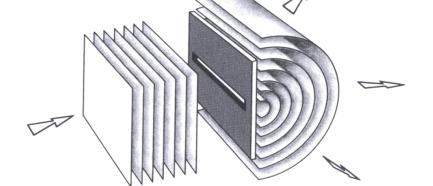
$$\nabla^2 = \frac{1}{r} \frac{\partial}{\partial r} \left( r \frac{\partial}{\partial r} \right) + \frac{1}{r^2} \frac{\partial^2}{\partial \theta^2} + \frac{\partial^2}{\partial z^2}$$

The requirement of cylindrical symmetry means that

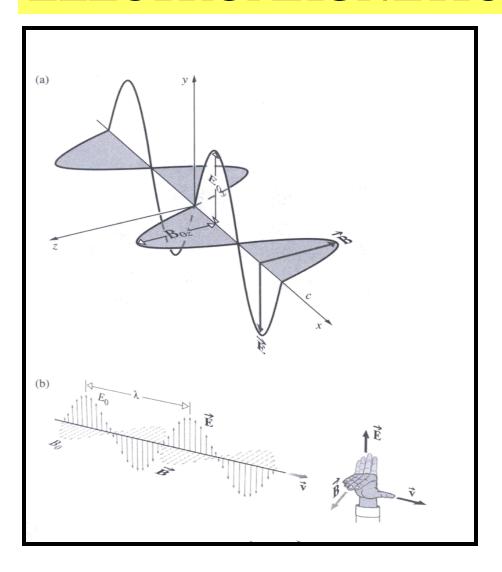
$$\psi(\vec{r}) = \psi(r, \theta, z) = \psi(r)$$

$$\psi(\mathbf{r},t) \approx \frac{A}{\sqrt{\rho}} e^{i(\mathbf{k}\rho \pm \omega t)}$$

cylindrical wavefunction



#### **ELECTROMAGNETIC WAVES**



$$E = E_o \sin(\vec{k}.\vec{r} - \omega t)$$

$$B = B_o \sin(\vec{k}.\vec{r} - \omega t)$$

$$E = cB$$

$$c = \frac{1}{\sqrt{\mu_0 \epsilon_0}}$$

# ENERGY IN AN EM WAVE

As with any wave, the EM wave transports energy

# **Energy density stored** in an E field

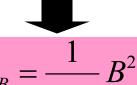


$$u_E = \frac{1}{2} \, \varepsilon_0 E^2$$

$$E = cB$$

$$c = 1/\sqrt{\mu_0 \varepsilon_0}$$

# **Energy density stored** in a B field



$$u_E = \frac{1}{2} \varepsilon_0 E^2 = \frac{1}{2} \varepsilon_0 c^2 B^2 = \frac{1}{2} \frac{\varepsilon_0}{\mu_0 \varepsilon_0} B^2 = \frac{1}{2 \mu_0} B^2 = u_B$$

The total energy density is shared between the constituent electric and magnetic fields

$$u = u_E + u_B$$

$$u = \varepsilon_0 E^2 = \frac{1}{\mu_o} B^2$$

# ENERGY IN AN EM WAVE

What is the EM energy flow associated with a traveling wave?

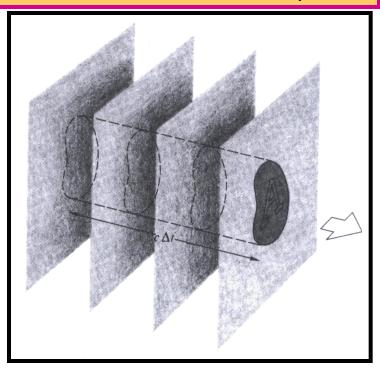
S: transport of energy per unit time across a unit area (W/m²)

During a very small interval  $\Delta t$ , only the energy contained in the cylindrical volume,  $u(c \Delta tA)$ , will cross A

$$S = \frac{uc\Delta tA}{\Delta tA} = uc = \frac{1}{\mu_o} EB$$

For isotropic media, the energy flows in the direction of the propagation of the wave:

$$\mathbf{S} = \frac{1}{\mu_o} \mathbf{E} \times \mathbf{B} = c^2 \varepsilon_o \mathbf{E} \times \mathbf{B}$$



**Poynting vector** 

# **IRRADIANCE**

Irradiance: AVERAGE energy per unit area per unit time

$$\mathbf{I} \equiv \left\langle \left| \mathbf{S} \right| \right\rangle_{\mathbf{T}} = \mathbf{c}^2 \mathbf{\epsilon}_{\mathbf{o}} \left| \mathbf{E}_{\mathbf{o}} \times \mathbf{B}_{\mathbf{o}} \right| \left\langle \sin^2(\mathbf{k.r} \pm \omega t) \right\rangle$$

$$I = \langle S \rangle_T = \frac{c\varepsilon_o}{2} E_o^2$$

$$I = c\varepsilon_o \langle E^2 \rangle_T = \frac{c}{\mu_o} \langle B^2 \rangle_T$$

**E** is considerably more effective at exerting forces and doing work on charges than **B** 

E is called the **OPTICAL FIELD** 

# LIGHT POLARIZATION

The direction of E is known as the polarization of the wave.

$$\vec{E} = E_o \sin(kz - \omega t) \hat{x}$$

$$\vec{\mathbf{B}} = \mathbf{B}_{o} \sin(k\mathbf{z} - \omega t) \hat{\mathbf{y}}$$

$$\vec{S} = \varepsilon_0 c E_0^2 \sin^2(kz - \omega t) \hat{z}$$

E is called the OPTICAL FIELD

The polaraization of an EM wave determines the direction of the force that the EM wave exerts on a charged particle in the path of the wave.

← Lorentz force law



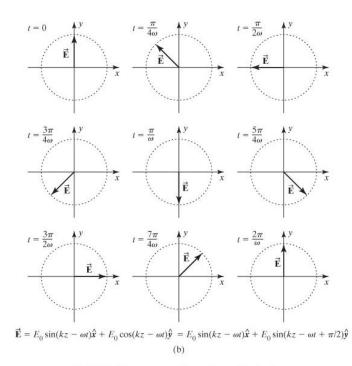
$$\vec{F} = Q(\vec{E} + \vec{v} \times \vec{B})$$

# LIGHT POLARIZATION

#### Linear polarization

# t = 0 $\vec{E} = 0$ $t = \frac{\pi}{4\omega}$ $t = \frac{\pi}{4\omega}$ $t = \frac{\pi}{4\omega}$ $t = \frac{\pi}{4\omega}$ $t = \frac{5\pi}{4\omega}$ $\vec{E} = 0$ $\vec{E} = 0$

#### circular polarization



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Random polarized, partial polarized light ← How to convert them to polarized light???

# DOPPLER EFFECT

$$\frac{\lambda'}{\lambda} = \sqrt{\frac{1 - \frac{\upsilon}{c}}{1 + \frac{\upsilon}{c}}}$$

υ is the relative velocity between the source and observer

- $v + \leftarrow$  approaching each other
- $\upsilon$   $\leftarrow$  separating from each other